

## Velocity of Sound in Nitrogen and Argon at High Pressures

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Measurements of the velocity of sound in N<sub>2</sub> and Ar are reported at the ice point and 294.26°K (70°F) and at pressures between 1 and 70 atm. The ratio of the specific heats as a function of pressure and temperature was calculated from the experimental results.

## INTRODUCTION

PRECISE measurements of the sound velocity in real gases at different temperatures and pressures offer a useful tool for studying the equation of state. Recently, Gyorog<sup>1</sup> developed a generalized equation of state derived from experimental measurements of compressibility factor and enthalpy and internal energy deviations. To gain insight into the validity of this equation of state for N<sub>2</sub> and Ar, velocity measurements were made in these gases at the ice point and

TABLE I. Sound velocity and ratio of specific heats in nitrogen.

| Temperature (°K) | Pressure (atm) | V <sub>0</sub> (exptl.) (m/sec) | V <sub>0</sub> (Gyorog) (m/sec) | γ      |
|------------------|----------------|---------------------------------|---------------------------------|--------|
| 273.15           | 1.00           | 337.04                          | 336.95                          | 1.4029 |
|                  | 10.00          | 338.17                          | 338.01                          | 1.4229 |
|                  | 30.00          | 341.37                          | 341.25                          | 1.4660 |
|                  | 50.00          | 345.88                          | 345.76                          | 1.5089 |
|                  | 70.00          | 351.55                          | 351.58                          | 1.5483 |
| 294.26           | 1.00           | 349.72                          | 349.76                          | 1.4014 |
|                  | 5.00           | 350.44                          | 350.37                          | 1.4097 |
|                  | 10.00          | 351.22                          | 351.20                          | 1.4185 |
|                  | 30.00          | 355.11                          | 355.12                          | 1.4552 |
|                  | 50.00          | 359.98                          | 360.09                          | 1.4900 |
|                  | 70.00          | 366.06                          | 366.10                          | 1.5246 |

294.26°K and at pressures from 1 to 70 atm. Also, knowledge of the molecular weight and the specific heat, along with the Gyorog equation, enabled velocity values to be calculated. The experimental data are compared with these calculated values, as well as with the data of van Itterbeek.<sup>2,3</sup>

## METHOD

The experimental method is based on the apparatus developed by the author and described in a previous paper.<sup>4</sup> The N<sub>2</sub> and Ar had purity of 99.95% or better.

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<sup>1</sup> D. A. Gyorog and E. F. Obert, *Am. Inst. Chem. Eng. J.* **10**, 625-631 (1964).

<sup>2</sup> A. van Itterbeek, W. DeRop, and G. Forrez, *Appl. Sci. Res.* **A6**, 5-6, 421-432 (1957).

<sup>3</sup> A. van Itterbeek, W. van Daël, and W. Grevendonk, *Physica* **25**, 7, 640-644 (1959).

<sup>4</sup> A. S. El-Hakeem, "A Refined Tube Method for Measuring the Sound Wavelength in Gases," *Am. J. Phys.* (to be published).

Temperature stability was achieved by running the tests in an air-conditioned room, the temperature of which did not change from 294.26°K by more than ±2 C° in a 24-h period. The entire sound tube was immersed in a water bath to ensure that no drastic changes in temperature (not more than ±0.1 C°) occurred during the test period. The ice-point temperature was controlled by immersing the sound tube in a bath of pure ice and distilled water.

Pressures of 1 and 5 atm were measured on a 130-in. manometer, while calibrated Heise absolute-pressure gauges measured the 10- to 70-atm points.

Extensive initial tests<sup>4</sup> showed that the apparatus was compatible with the Helmholtz-Kirchoff (H-K) equation. Thus the H-K constant for the tube (β<sub>T</sub>) was equal, within experimental accuracy, to the theoretical value (β) and, therefore, could be calculated. The velocity of sound in the tube was then measured at one or more frequencies and corrected with the calculated β to yield the Laplacian free-gas velocity V<sub>0</sub>.

The sound velocities were also calculated with the aid of Eq. (1), using the generalized virial coefficients of Gyorog<sup>1</sup> (or the generalized virial coefficients derived from the Lennard-Jones potential)<sup>5,6</sup>:

$$V_0 = \left\{ (RT/M) [\phi + (\Psi/x)(\gamma^0 - 1)] \right\}^{1/2}, \quad (1)$$

where

$$\begin{aligned} \phi &= 1 + 2B^*\rho^* + 3C^*\rho^{*2} + 4D^*\rho^{*3}, \\ \Psi &= \left[ 1 + \left( B^* + T^* \frac{dB^*}{dT^*} \right) \rho^* + \left( C^* + T^* \frac{dC^*}{dT^*} \right) \rho^{*2} \right. \\ &\quad \left. + \left( D^* + T^* \frac{dD^*}{dT^*} \right) \rho^{*3} \right]^2, \\ x &= 1 - (\gamma^0 - 1) \left[ \left( 2T^* \frac{dB^*}{dT^*} + T^{*2} \frac{d^2B^*}{dT^{*2}} \right) \rho^* \right. \\ &\quad \left. + \left( T^* \frac{dC^*}{dT^*} + \frac{1}{2} T^{*2} \frac{d^2C^*}{dT^{*2}} \right) \rho^{*2} + \left( \frac{2}{3} T^* \frac{dD^*}{dT^*} + \frac{1}{3} T^{*2} \frac{d^2D^*}{dT^{*2}} \right) \rho^{*3} \right], \end{aligned}$$

γ<sup>0</sup> is the ratio of heat capacities for the ideal gas state;

<sup>5</sup> J. O. Hirschfelder, C. F. Curtiss, and R. B. Bird, *Molecular Theory of Gases and Liquids* (John Wiley & Sons, Inc., New York, 1954).

<sup>6</sup> For the range of the reduced temperatures covered in this paper, the sound velocity values calculated with these virial coefficients are essentially identical to those calculated from Gyorog's equation.

TABLE I  
Temperature  
(°K)

273.15

294.26

R, T, M are  
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The ratio  
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FIG. 1. γ (X, present data<sup>2</sup>.)

<sup>7</sup> V. Hov  
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